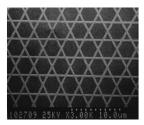
Hierarchy of critical behaviours in the interacting three colouring model on the honeycomb lattice

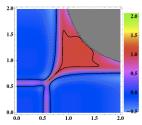
Claudio Castelnovo

Hubbard Theory Consortium, Department of Physics, Royal Holloway University



Pierre Pujol LPT IRSAMC Université Paul Sabatier

Jacob Simmons
JFI, University of Chicago



Open Statistical Physics, The Open University, Milton Keynes UK, March 2, 2011

Outline

- ▶ (semi) historical introduction to the three-colouring model
- ▶ the role of interactions: a "continuously varying" central charge in a unitary system?
- further insight introducing colour-dependent coupling constants
- conclusions and open questions

The model Baxter 1970

JOURNAL OF MATHEMATICAL PHYSICS VOLUME 11, NUMBER 3 MARCH 1970

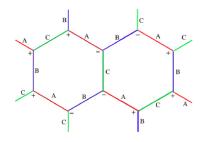
Colorings of a Hexagonal Lattice

R. J. BAXTER

Mathematics Department, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139

(Received 13 May 1969)

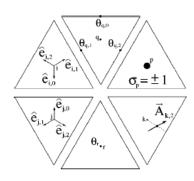
The number of ways W^L of coloring the bonds of a hexagonal lattice of L sites (L large) with three colors so that no adjacent bonds are colored alike is calculated exactly, giving $W = 1.20872 \cdots$. This is equivalent to counting the number of 4-colorings of the faces of the lattice and can also be regarded as a multiple-dimer problem. If one introduces activities corresponding to certain vertex configurations, then the system is found to have an infinite-order phase transition between two ordered states.



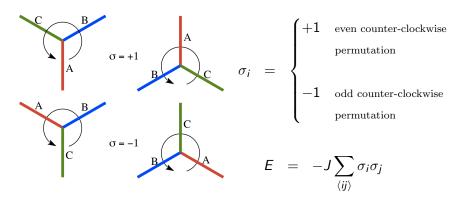
- hard constraint: no two bonds with equal colour can meet at any vertex
- extensive degeneracy
- power-law correlations $SU(3)_{k=1}$ WZNW CFT N. Read



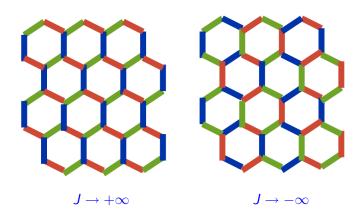
Higgins et al. 2000



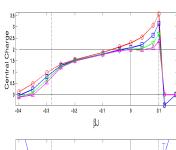
Castelnovo et al. 2004

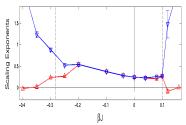


- ▶ glassy dynamics (freezing into polycrystalline metastable states) [Phys. Rev. B 69, 104529 (2004)]
- ► unusual <u>critical behaviour</u>



- $ightharpoonup J > 0 \longrightarrow$ favours parallel bonds having the same colour
- ▶ $J < 0 \longrightarrow$ favours alternating sequences of two colours around hexagonal plaquettes





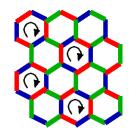
- ► $J = 0 \longrightarrow SU(3)_{k=1}$ Wess-Zumino-Novikov-Witten conformal field theory (central charge c = 2)
- J > 0 → line of critical points ending in a first order transition (approx. constant c and scaling exponents)
- ► J < 0 → continuously varying central charge?

Kolmogorov 'c' theorem (for a unitary CFT) forbids a continuously varying central charge!

new insight from colour-dependent couplings (A=red, B=blue, C=green):

$$E = -\sum_{\langle ij \rangle} J_{\alpha_{ij}} \sigma_i \sigma_j$$
 $\alpha_{ij} = A, B, C$

- ▶ $J_A = J_B = -\infty$ lsing model on the triangular lattice with nn interaction J_C
- $\blacktriangleright J_A = -\infty, J_B = J_C \equiv J \longrightarrow \text{dimer}$ model on the honeycomb lattice with loop tension J Jacobsen+Alet 2009

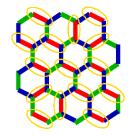


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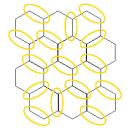


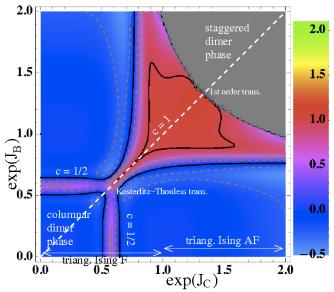
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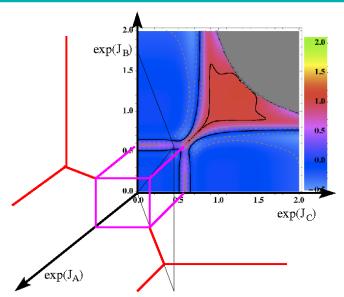
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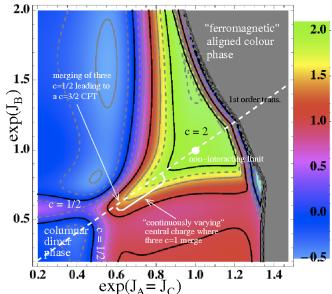
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Conclusions and open questions

- ► competition between short-range interaction and hard-constraints ⇒ wide range of critical behaviours
- new insight on how to interpret the different critical behaviours comes from colour-dependent couplings:
 - two c = 1/2 (free majoranas) merge into a c = 1 (free boson)
 - three c=1/2 (free majoranas) merge into a c=3/2 as if they were independent (!) [supersymmetric point?]
 - two c = 1 (free bosons) merge into a c = 2 (free boson)
 - ▶ three c = 1 (free bosons) merge into $c \in (3/2, 2)$



the paradox of the "continuously varying" central charge remains unsolved

